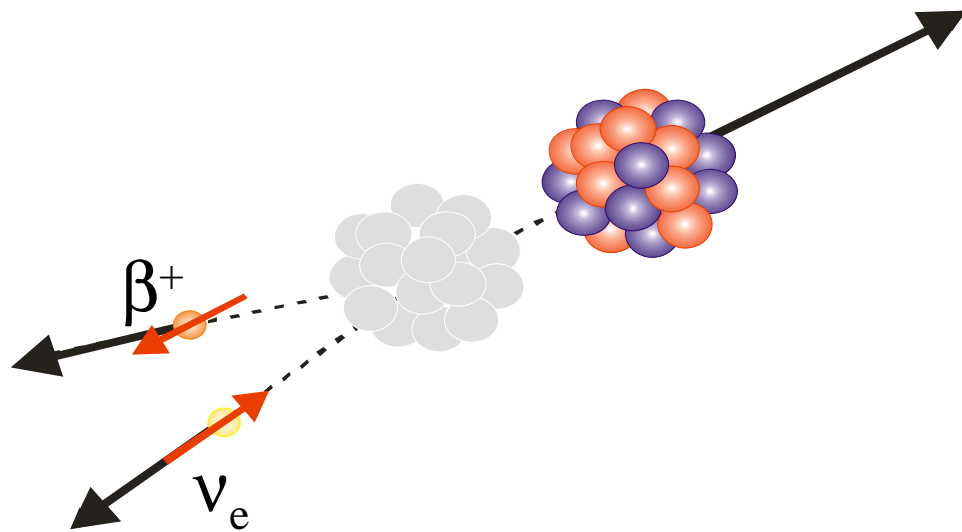


Fundamental interactions experiments with polarized trapped nuclei



DESIR meeting
Leuven, 26-28 May 2010

Nathal Severijns
Kath. University Leuven, Belgium

1. searches for exotic weak currents (non V-A)

- tensor currents
- scalar currents

2. symmetry tests

- parity
- time reversal / CP violation

3. determine V_{ud} and test CVC

distribution in

- electron and neutrino directions and in
- electron energy

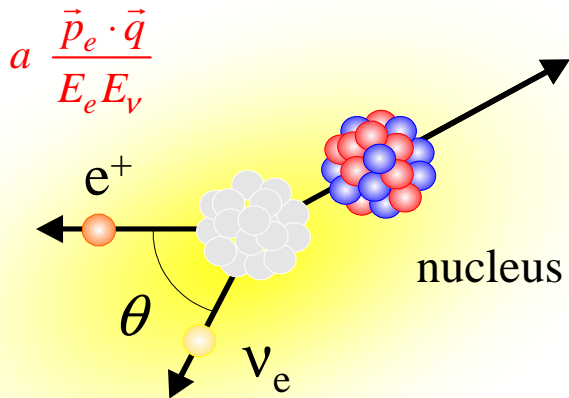
from polarized nuclei :

Correlation measurements

$$\omega(\langle \vec{J} \rangle | E_e, \Omega_e, \Omega_\nu) dE_e d\Omega_e d\Omega_\nu$$

$$\propto \frac{F(\pm Z, E_e)}{\text{Fermi function}} \frac{p_e E_e (E_0 - E_e)^2 dE_e d\Omega_e d\Omega_\nu}{\text{phase space}}$$

$$\times \zeta \left\{ 1 + a \frac{\vec{p}_e \cdot \vec{q}}{E_e E_\nu} + b \frac{\gamma m_e}{E_e} + A \frac{\vec{J} \cdot \vec{p}_e}{J E_e} + R \vec{\sigma} \cdot \frac{\vec{J}}{J} \times \frac{\vec{p}_e}{E_e} + \dots \right\}$$



$a \frac{\vec{p}_e \cdot \vec{q}}{E_e E_\nu}$
 β -v correlation

$b \frac{\gamma m_e}{E_e}$
 Fierz interference term
 ($b \equiv 0$ in standard model)

$A \frac{\vec{J} \cdot \vec{p}_e}{J E_e}$
 β -asymmetry

$R \vec{\sigma} \cdot \frac{\vec{J}}{J} \times \frac{\vec{p}_e}{E_e}$
 R-correlation

$$\tilde{X} = \frac{X}{1 + b \frac{\gamma m_e}{E_e}}$$

J,D, Jackson, S.B. Treiman, H.W. Wyld, Nucl. Phys. 4 (1957) 206

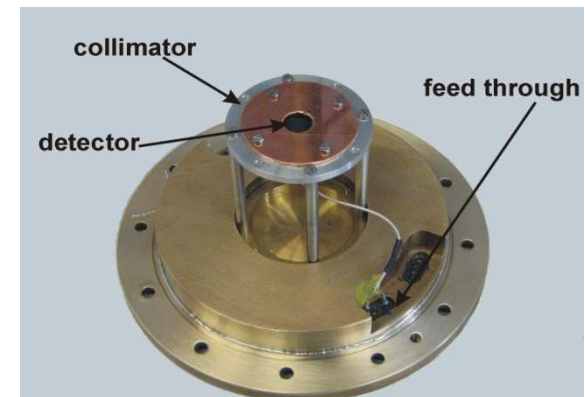
^{114}In and ^{60}Co beta asymmetry parameter, A

Wauters et al.



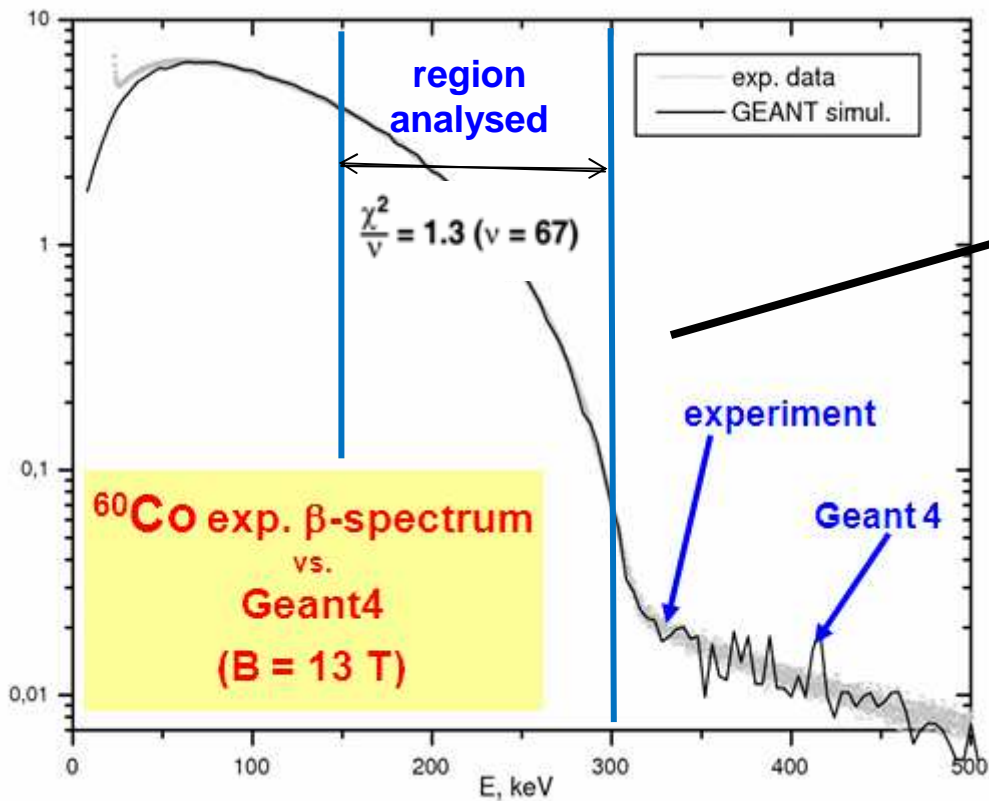
in 13 T external polarizing field;
used GEANT4 code to account for :

- detection geometry
- magnetic field effects
- scattering



Si p-I-n diode
(500 μm , $\varnothing = 9 \text{ mm}$)
operating at 10 K

Leuven $^3\text{He} - ^4\text{He}$ dilution refrigerator setup



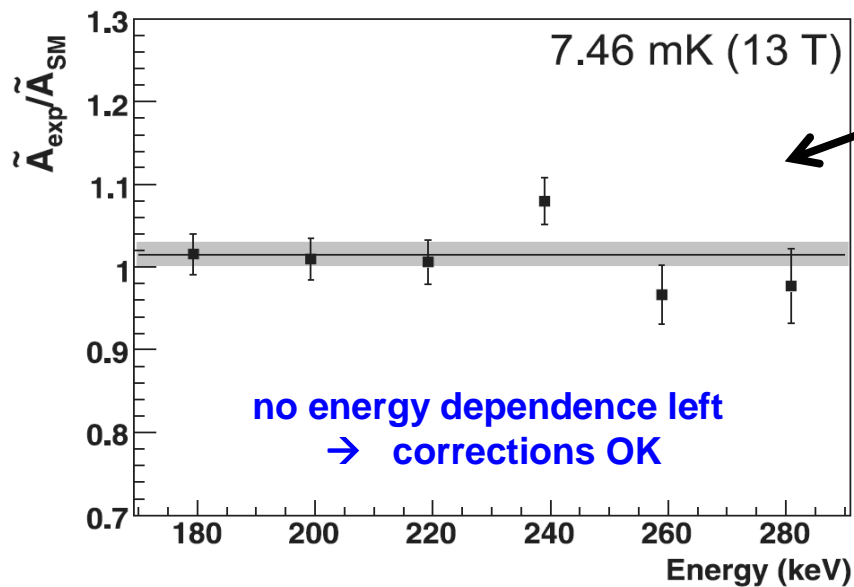
$$W(\theta) = \frac{N(\theta)_{\text{pol}}}{N(\theta)_{\text{unpol}}} = 1 + \tilde{A} P \left(\frac{v}{c} Q \cos\theta \right)$$

(P from anisotropy of γ -rays)

Geant 4

Analysis of measured and simulated spectra:

$$\frac{[W(\theta) - 1]_{\text{exp}}}{[W(\theta) - 1]_{\text{Geant}}} = \frac{\left[\tilde{A}_{\text{exp}} P \frac{v}{c} Q \cos\theta \right]_{\text{exp}}}{\left[\tilde{A}_{\text{SM}} P \frac{v}{c} Q \cos\theta \right]_{\text{Geant}}}$$



$$A_{\text{exp}}(^{60}\text{Co}) = -1.014 (12)_{\text{stat}} (16)_{\text{syst}}$$

$$(A_{\text{SM}} = -0.987(9))$$

F. Wauters et al., submitted

$$A_{\text{exp}}(^{114}\text{In}) = -0.994 (10)_{\text{stat}} (10)_{\text{syst}}$$

$$(A_{\text{SM}} = -1.000)$$

(most precise result for A_{nuclear} ever !)

F. Wauters et al., Phys. Rev. C 80 (2009) 062501(R)

(^{67}Cu in progress)

(M)LRS-models

$$W_1 = W_L \cos\zeta - W_R \sin\zeta$$

$$W_2 = W_L \sin\zeta + W_R \cos\zeta$$

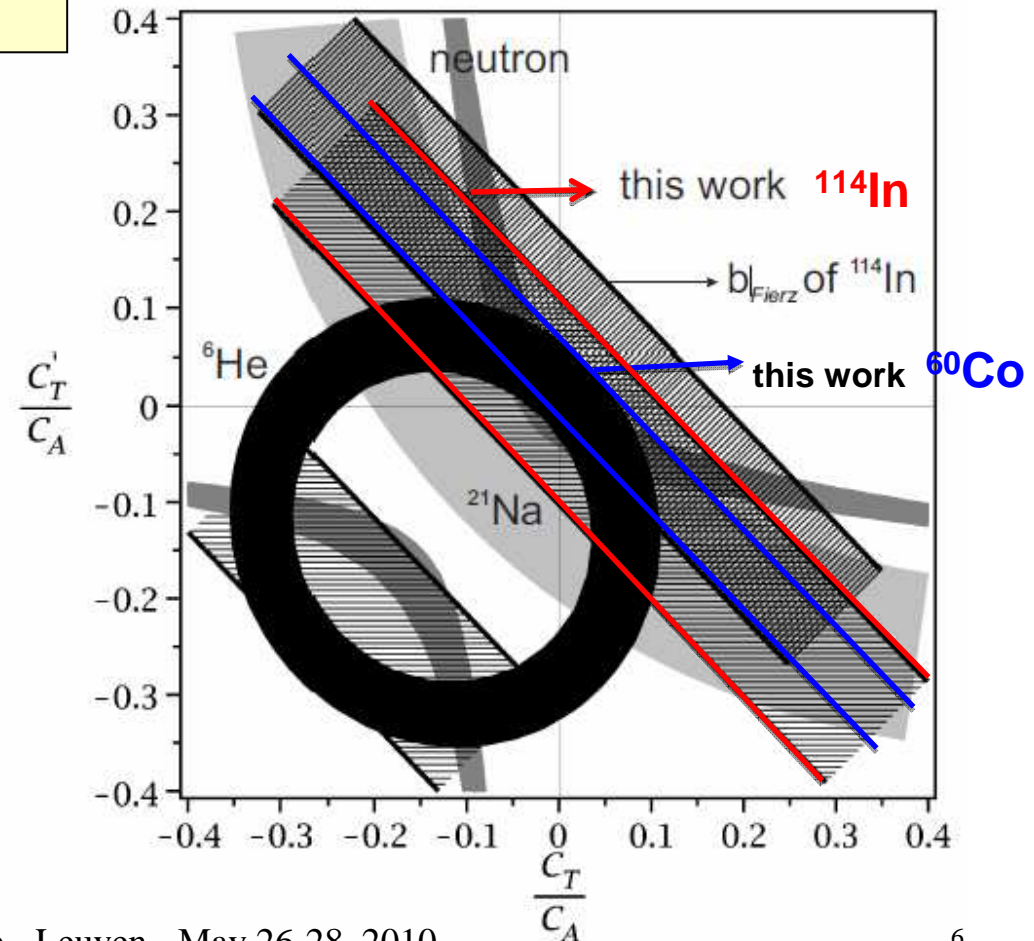
$$\delta = m_1^2 / m_2^2$$

$$^{114}\text{In} : M(W_2) > 230 \text{ GeV}/c^2 \quad (90\% \text{ C.L.})$$

$$^{60}\text{Co} : M(W_2) > 245 \text{ GeV}/c^2 \quad (90\% \text{ C.L.})$$

major systematic errors:

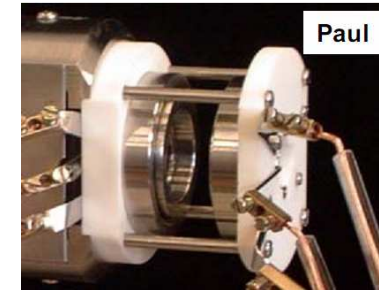
- performance of GEANT code (scattering)
- determination of nuclear polarization



Polarizing atoms/ions in a particle trap:

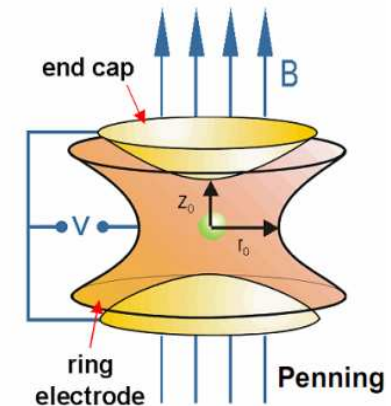
- **Paul trap** :
(LPC-GANIL)

optical pumping of ion cloud
in magnetic holding field



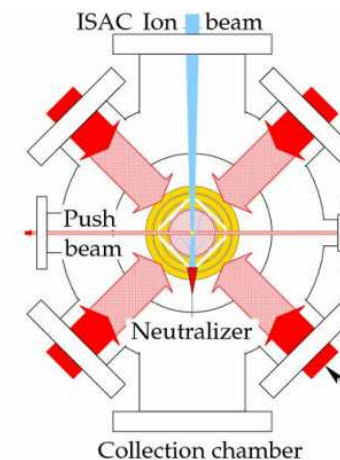
- **Penning trap** :
(WITCH-ISOLDE, DESIR)

collinear polarization by optical pumping
in beam line before trap



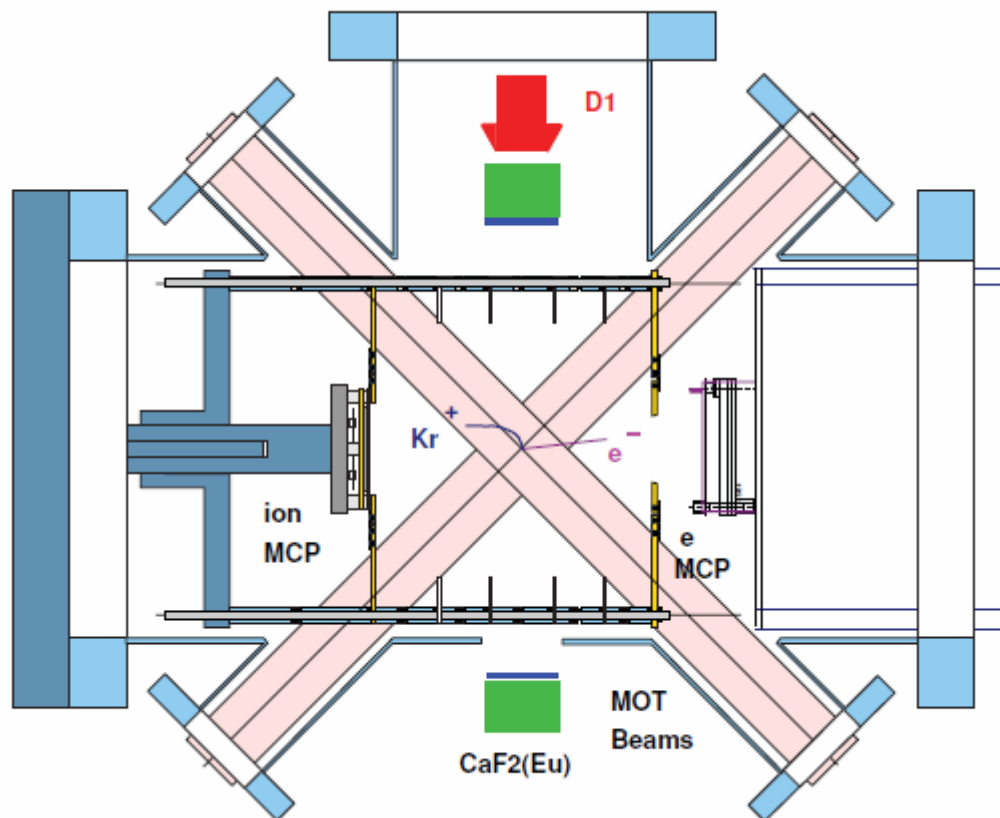
- **MOT trap** :
(TRIUMF, Berkeley, KVI)

optical pumping of ion cloud
in magnetic holding field



Alternative : stop ions in superfluid He and
polarize by optical pumping (cf. talk by T. Shimoda)

Example of polarized atoms in MOT: neutrino asymmetry parameter for ^{37}K



TRINAT MOT trap @ TRIUMF

$$B_\nu(^{37}\text{K}) = -0.755(24)$$

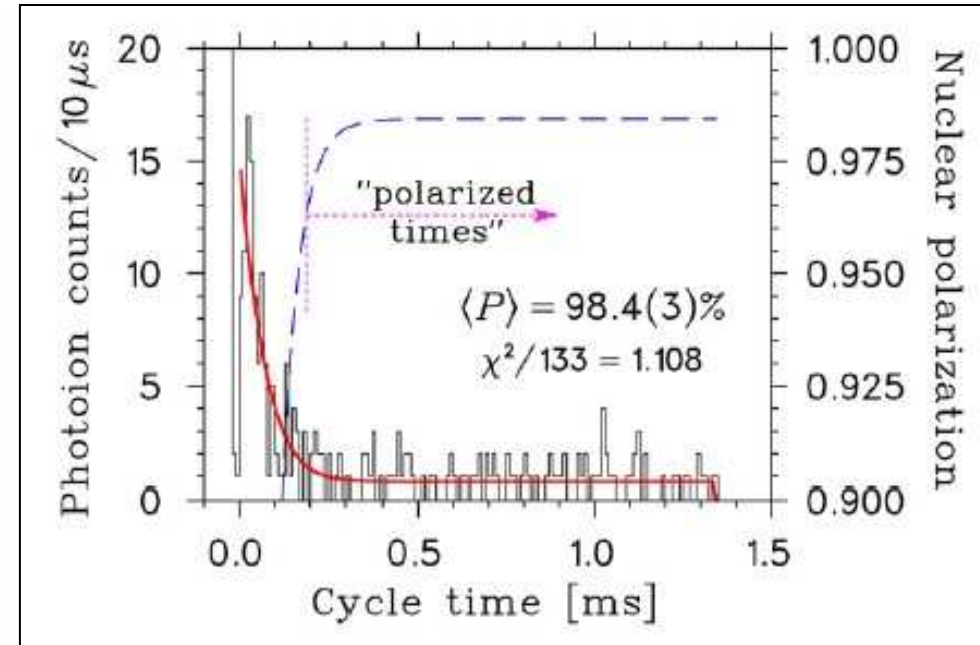
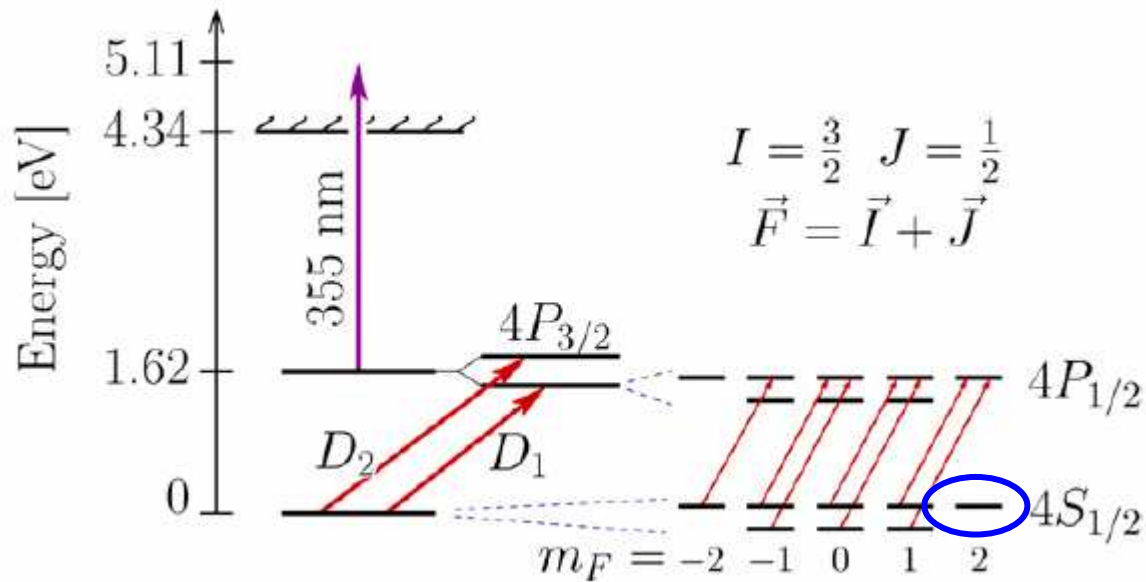
$$B_\nu^{\text{SM}} = -0.779(6)$$

5 % precision
very difficult to determine
nuclear polarization precisely !

$$M_{W2} > 180 \text{ GeV}/c^2 \text{ (90 \% C.L.)}$$

^{37}K – D. Melconian, J.A. Behr et al., Phys. Lett. B 649 (2007) 370

Polarization by **optical pumping** and determination of nuclear polarization via **photoionization** in a MOT



^{37}K

$\langle P_{\sigma+} \rangle = (+97.7 \pm 0.4_{-0.5}^{+0.2})\%$

$\langle P_{\sigma-} \rangle = (-95.8 \pm 1.0_{+1.3}^{-0.4})\%$

D. Melconian, J.A. Behr et al., Phys. Lett. B 649 (2007) 370

^{80}Rb

$P = 0.53 \pm 0.03$

J.R.A. Pitcairn, J.A. Behr et al., Phys. Rev. C 79 (2009) 015501

Longitudinal polarization of positrons emitted by polarized nuclei

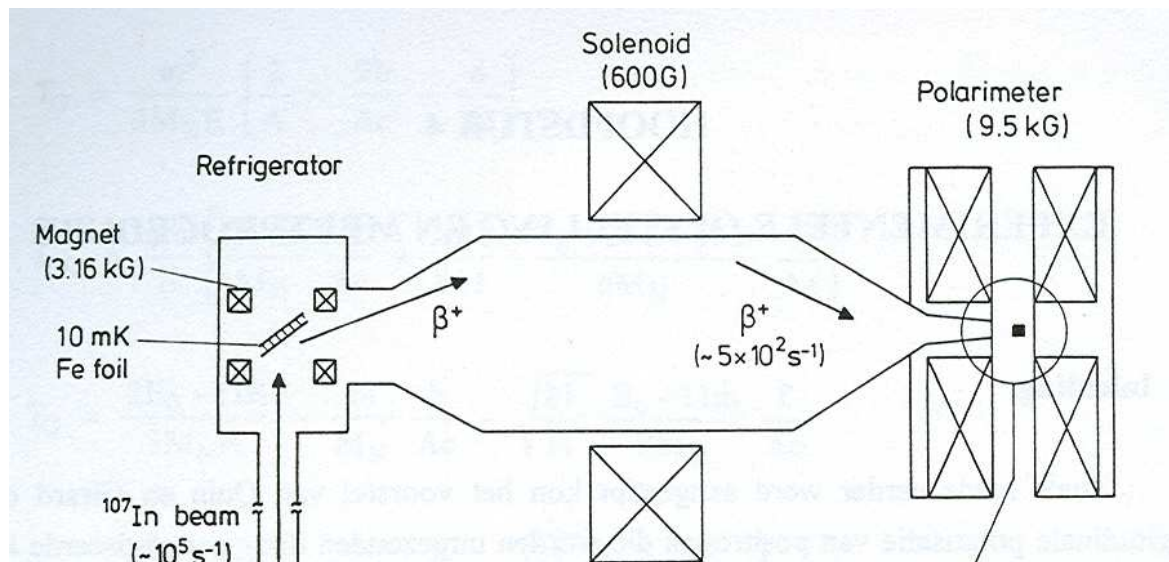
experimental quantity

$$R \equiv P^- / P^+ = R_{SM} [1 + k \Delta]$$

$$\Delta = (\delta + \zeta)^2$$

with :

P^- (P^+) is β particle longitudinal polarization for β 's emitted opposite to (in the direction of) the polarized nuclear spin vector (P^0 for unpolarized nuclei)

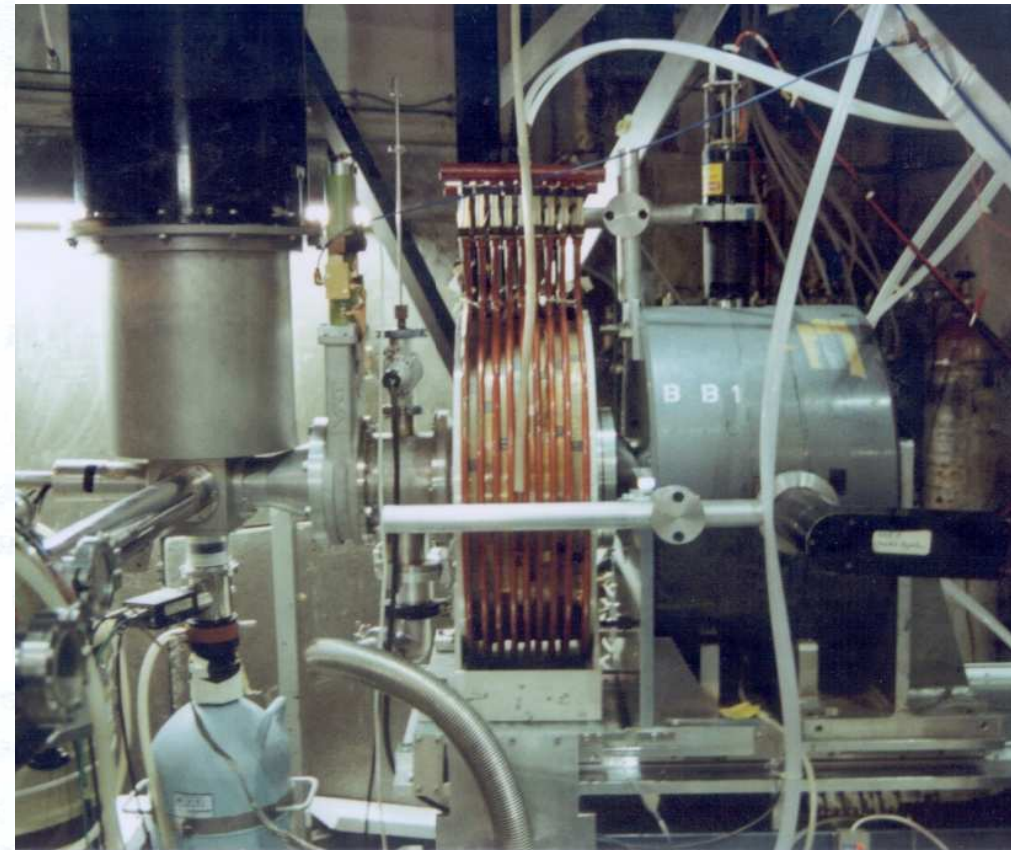


(M)LRS-models

$$W_1 = W_L \cos \zeta - W_R \sin \zeta$$

$$W_2 = W_L \sin \zeta + W_R \cos \zeta$$

$$\delta = (M_{W1})^2 / (M_{W2})^2$$

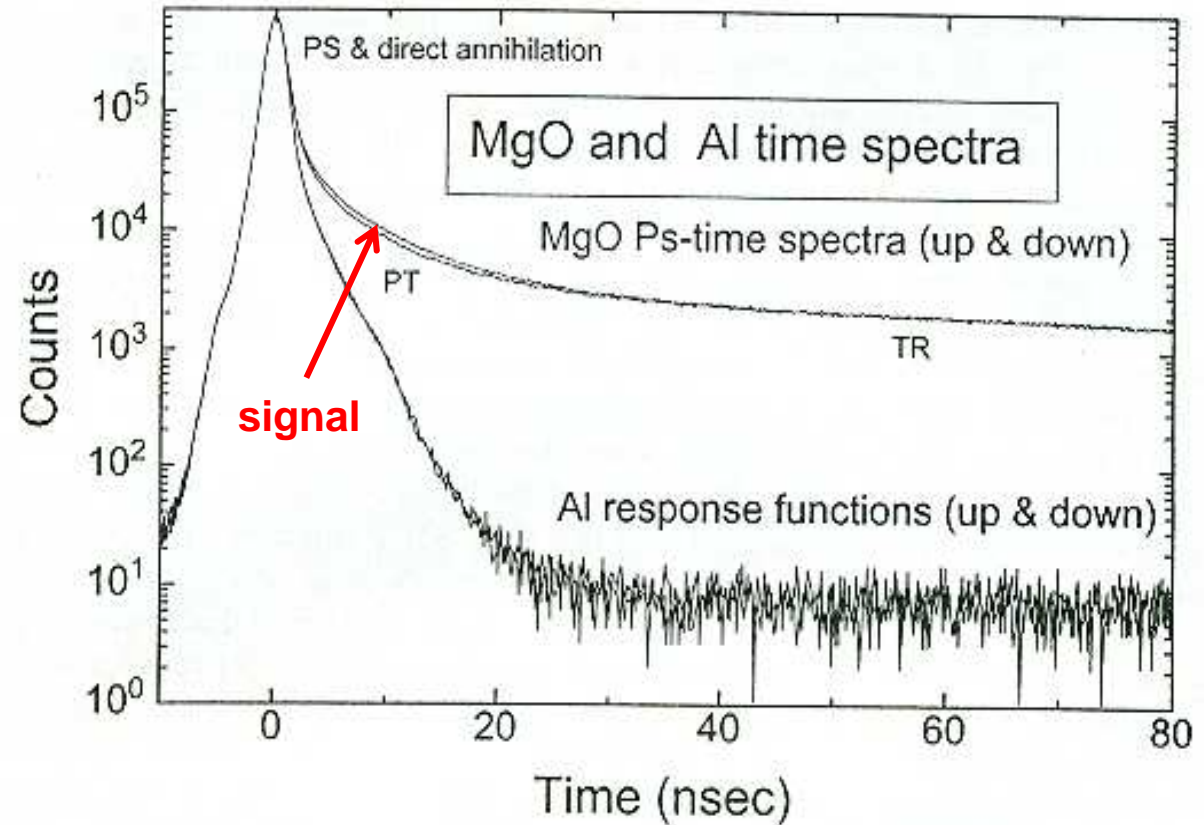


Results :

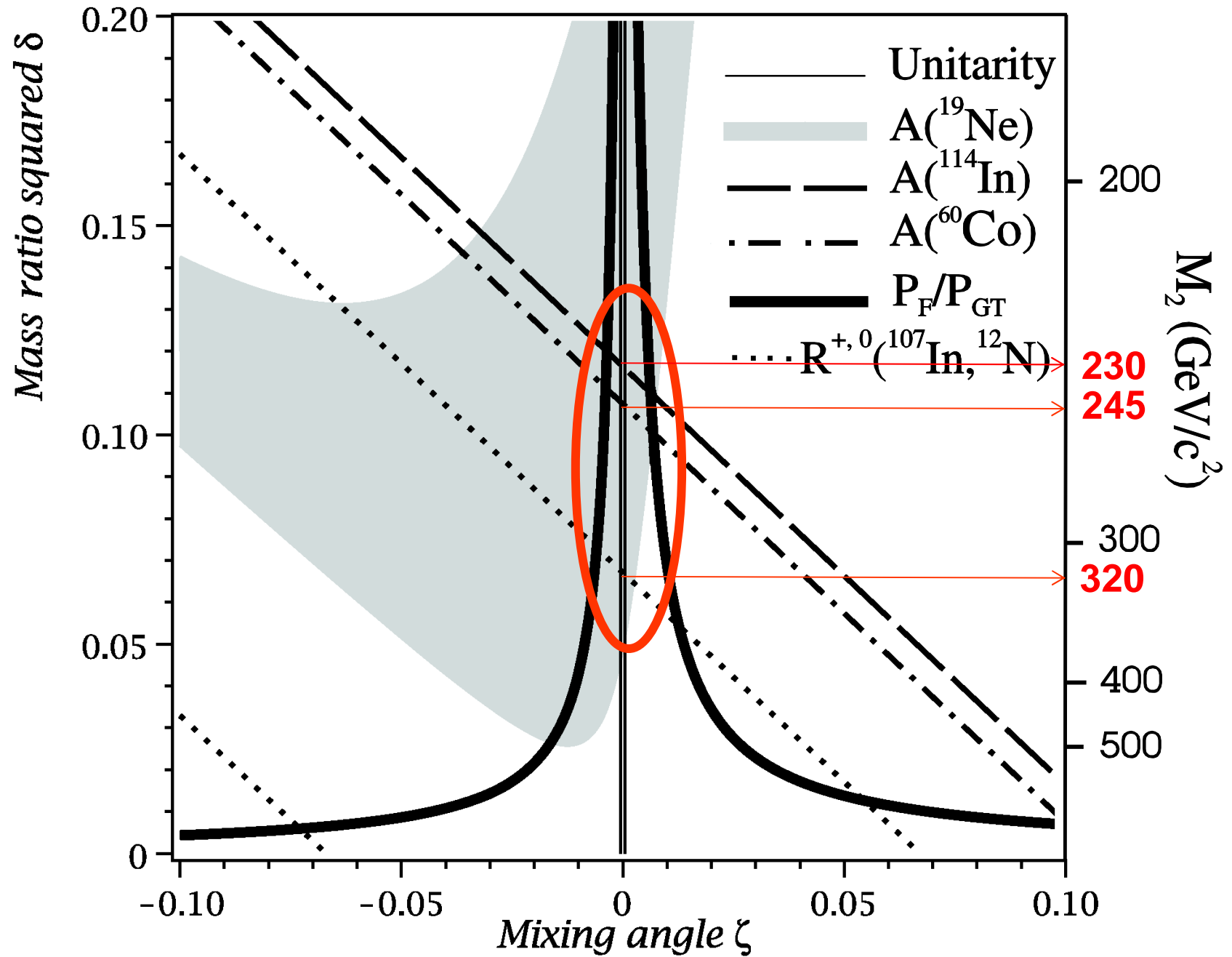
observable nucleus $(\delta + \zeta)^2$

P ⁻ / P ⁺	¹⁰⁷ In	-0.0003(58)	1)
P ⁻ / P ⁺	¹² N	0.0064(76)	2)
P ⁻ / P ⁺	¹² N	-0.0001(34)	3)
P ⁻ / P ⁰	¹⁰⁷ In	0.0021(17)	4)
average		0.0017(14)	

$M_{W_2} > 320 \text{ GeV}/c^2$ (90 % C.L.)



- 1) N. Severijns et al., PRL 70 (1993) 4047, 73 (1994) 611
- 2) M. Allet et al., Phys. Lett. B363 (1996) 139
- 3) E. Thomas et al., Nucl. Phys. A694 (2001) 559
- 4) N. Severijns et al., Nucl. Phys. A629 (1998) 423c



3. Relative positron polarization measurements with polarized nuclei

$$R \equiv \frac{P^-}{P^+} = R_0 \left[1 - \frac{8\beta^2 A_{\text{exp}}}{\beta^4 - A_{\text{exp}}^2} (\delta + \zeta)^2 \right] \quad \text{with} \quad R_0 = \frac{[\beta^2 - A_{\text{exp}}][1 + A_{\text{exp}}]}{[\beta^2 + A_{\text{exp}}][1 - A_{\text{exp}}]}$$

sensitivity factor k
 (>> for $\beta^2 = A_{\text{exp}}$!)

$$A_{\text{exp}} = \beta P A \langle \cos \theta \rangle$$

$\beta = v/c$; $P = \text{nuclear polarization}$

Isotope	k	sensitivity to M_{W2} [GeV/c ²] (90% C.L.)
¹⁷ F	32.2	604
²¹ Na	12.5	477
²⁵ Al	15.7	505
⁴¹ Sc	18.8	528
¹² N	17.9	522
¹⁰⁷ In	25.7	571


significant potential compared to present lower limit of 320 GeV/c²

(for $P = 0.80$, and a 1% precision on R/R_0)

T = 1/2 superallowed mirror transitions

$$\mathcal{F}t^{\text{mirror}} \equiv f_V t (1 + \delta'_R) (1 + \delta_{\text{NS}}^V - \delta_C^V) = \frac{2\mathcal{F}t^{0^+ \rightarrow 0^+}}{\left(1 + \frac{f_A}{f_V} \rho^2\right)} \quad \text{with } \rho = C_A M_{GT} / C_V M_F$$

Parent nucleus	$\mathcal{F}t$ (s)	$\delta\mathcal{F}t$ (%)
^3H	1135.3 ± 1.5	0.13
^{11}C	3933 ± 16	0.41
^{13}N	4682.0 ± 4.9	0.10
^{15}O	4402 ± 11	0.25
^{17}F	2300.4 ± 6.2	0.27
^{19}Ne	1718.4 ± 3.2	0.19
^{21}Na	4085 ± 12	0.29
^{23}Mg	4725 ± 17	0.36
^{25}Al	3721.1 ± 7.0	0.19
^{27}Si	4160 ± 20	0.48
^{29}P	4809 ± 19	0.40
^{31}S	4828 ± 33	0.68
^{33}Cl	5618 ± 13	0.23
^{35}Ar	5688.6 ± 7.2	0.13
^{37}K	4562 ± 28	0.61
^{39}Ca	4315 ± 16	0.37
^{41}Sc	2849 ± 11	0.39
^{43}Ti	3701 ± 56	1.51
^{45}V	4382 ± 99	2.26



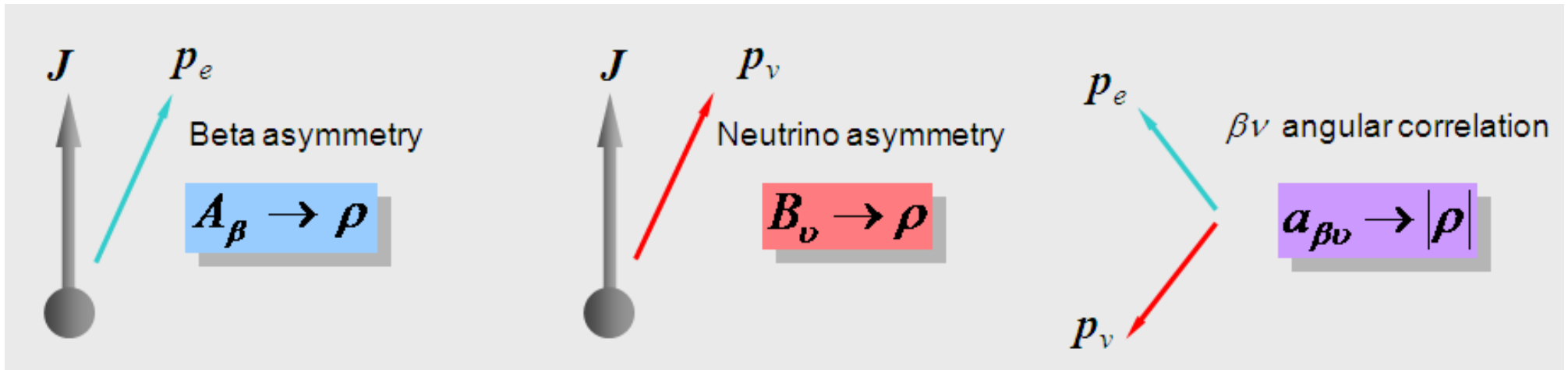
$$Ft_0 = Ft^{\text{mirror}} \left(1 + \frac{f_A}{f_V} \rho^2 \right) = 2Ft^{0^+ \rightarrow 0^+}$$

$$= \frac{K}{G_F^2 V_{ud}^2 (1 + \Delta_R^V)}$$

← accuracy of 0.1 % to 0.4 % for most cases

[NS, I.S. Towner et al., Phys. Rev. 78 (2008) 055501]

- extract mixing ratio $\rho = C_A M_{GT} / C_V M_F$ from correlation measurements:

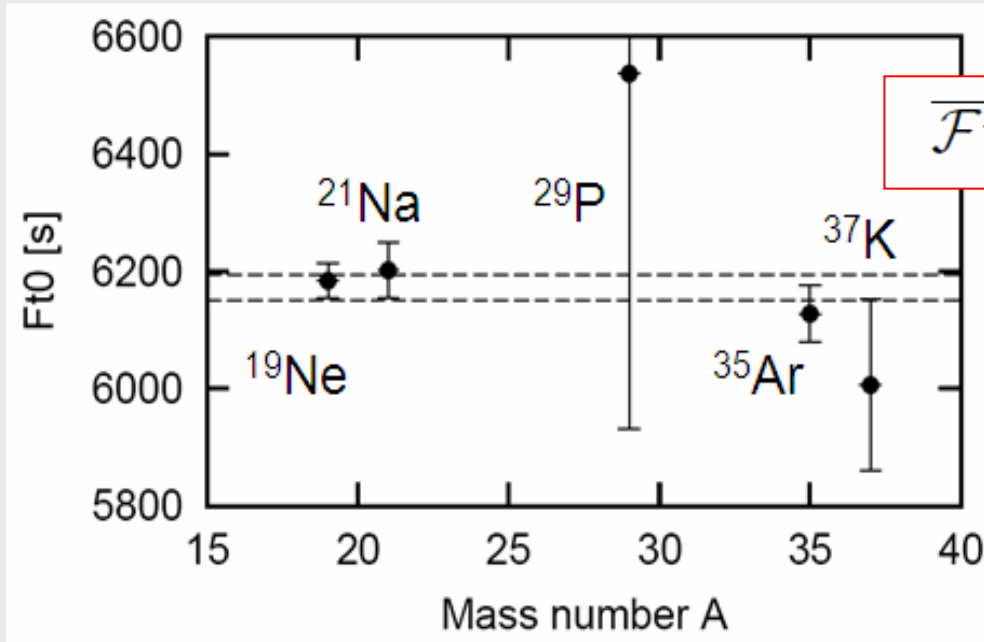


- there are 35 candidates between ^3H and ^{83}Mo , near the $N = Z$ line
(best are the ones with $A < 45$ about)
- correlation measurements have been carried out for:

^{17}F , ^{19}Ne , ^{21}Na , ^{29}P , ^{35}Ar and ^{37}K

Is the strength of the vector coupling the same in all $T=1/2$ transitions ?

[O. Naviliat-Cuncic & N.S., Phys. Rev. Lett. 102 (2009) 142302



$$\overline{Ft}_0 = 6173 \pm 22 \text{ s}$$

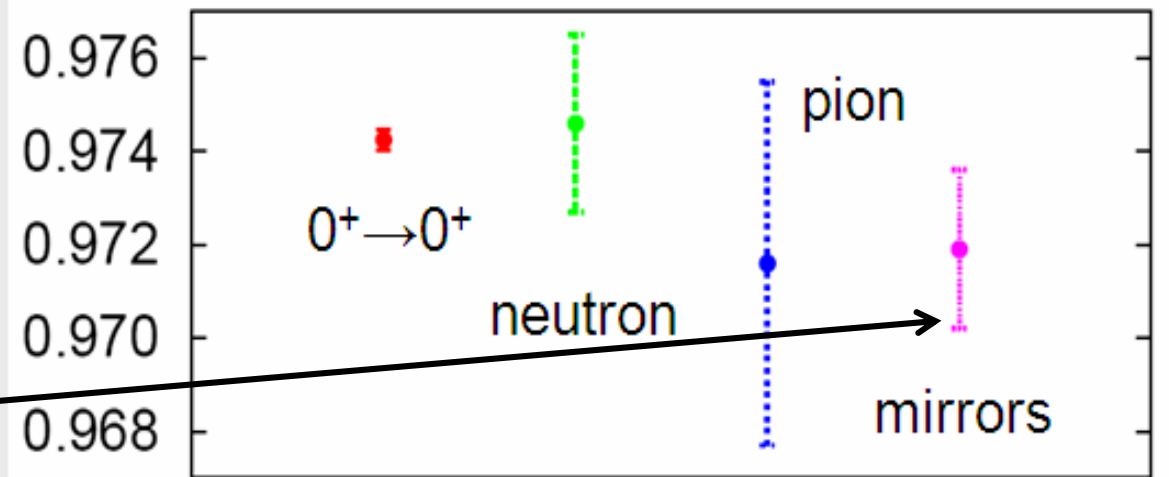
$$Ft_0 = 2Ft(0^+ \rightarrow 0^+)$$

- First consistent test of CVC from a set of nuclear transitions other than super-allowed pure Fermi

- New value of $|V_{ud}|$

$$|V_{ud}| = 0.9719(17)$$

reduce by new measurements of Ft values and correlation coefficients



- sensitivity of a and A for mirror nuclei, best cases :

Isotope	$\beta\nu$ correlation: $\Delta a = 0.5\%$		β asymmetry parameter $\Delta A = 0.5\%$	
	ΔV_{ud} (present Ft-value)	ΔV_{ud} (no error from Ft)	ΔV_{ud} (present Ft-value)	ΔV_{ud} (no error from Ft)
^3H	0.0011	0.0010	0.0011	0.0009
^{13}N	0.0017	0.0017	-	-
^{15}O	0.0020	0.0016	-	0.0020
^{17}F	0.0018	0.0014	-	-
^{19}Ne	0.0014	0.0010	0.0014	0.0011
^{25}Al	0.0020	0.0018	-	-
^{29}P	-	0.0018	0.0024	0.0014
^{33}Cl	0.0021	0.0018	0.0013	0.0006
^{35}Ar	0.0019	0.0018	0.0007	0.0004

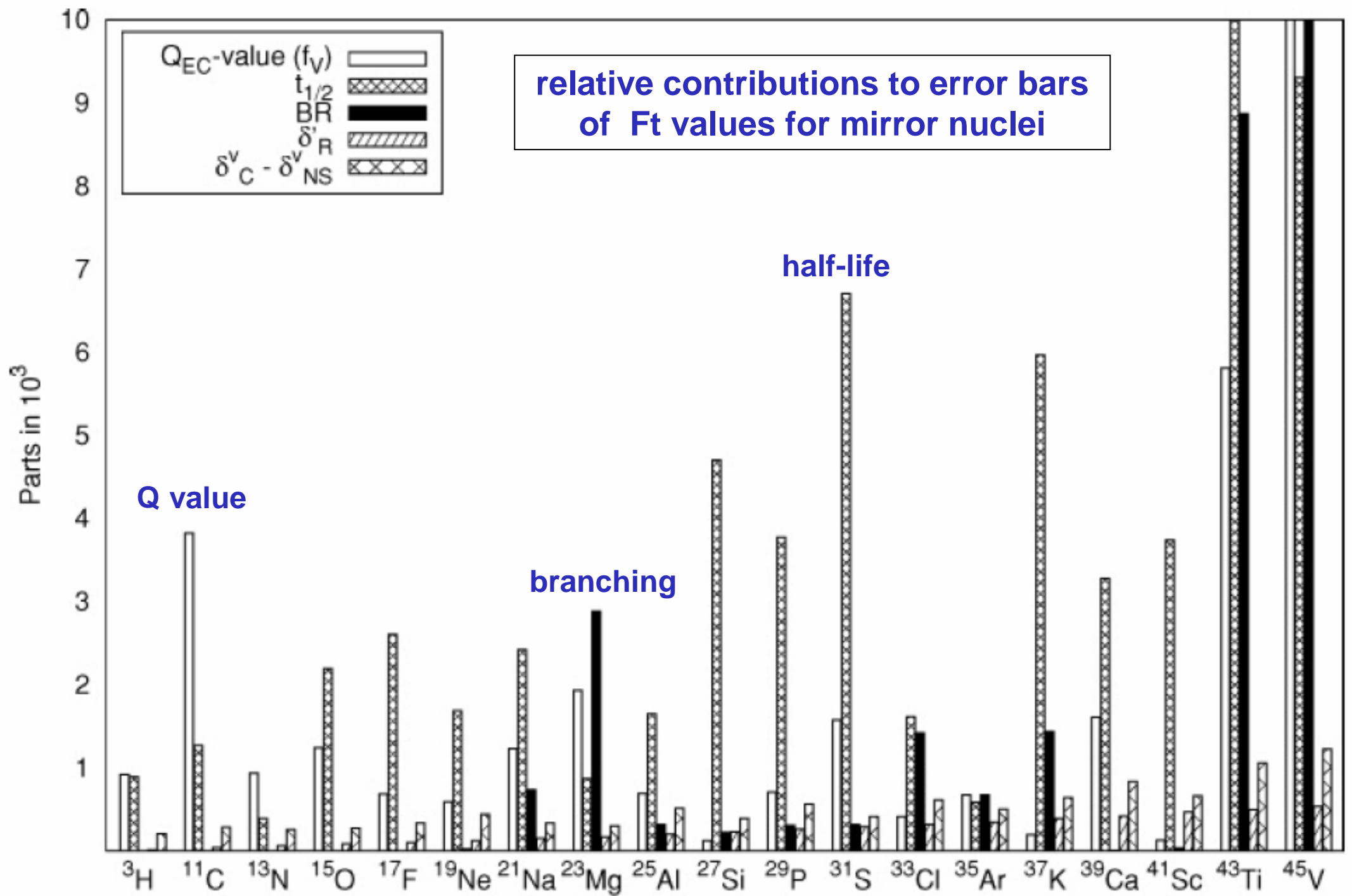
requires $\delta(\text{Ft}) \sim 5 \times 10^{-4}$
i.e.
factor 2 -10
better

Note: - ΔV_{ud} from $a_{\beta\nu}$ for all mirror transitions up to $^{39}\text{Ca} \leq 0.0018$ if no error from Ft
- $|V_{ud}| (0^+ \rightarrow 0^+) = 0.97425 \pm 0.00022$ and $|V_{ud}| (\text{mirror}) = 0.9719 \pm 0.0017$

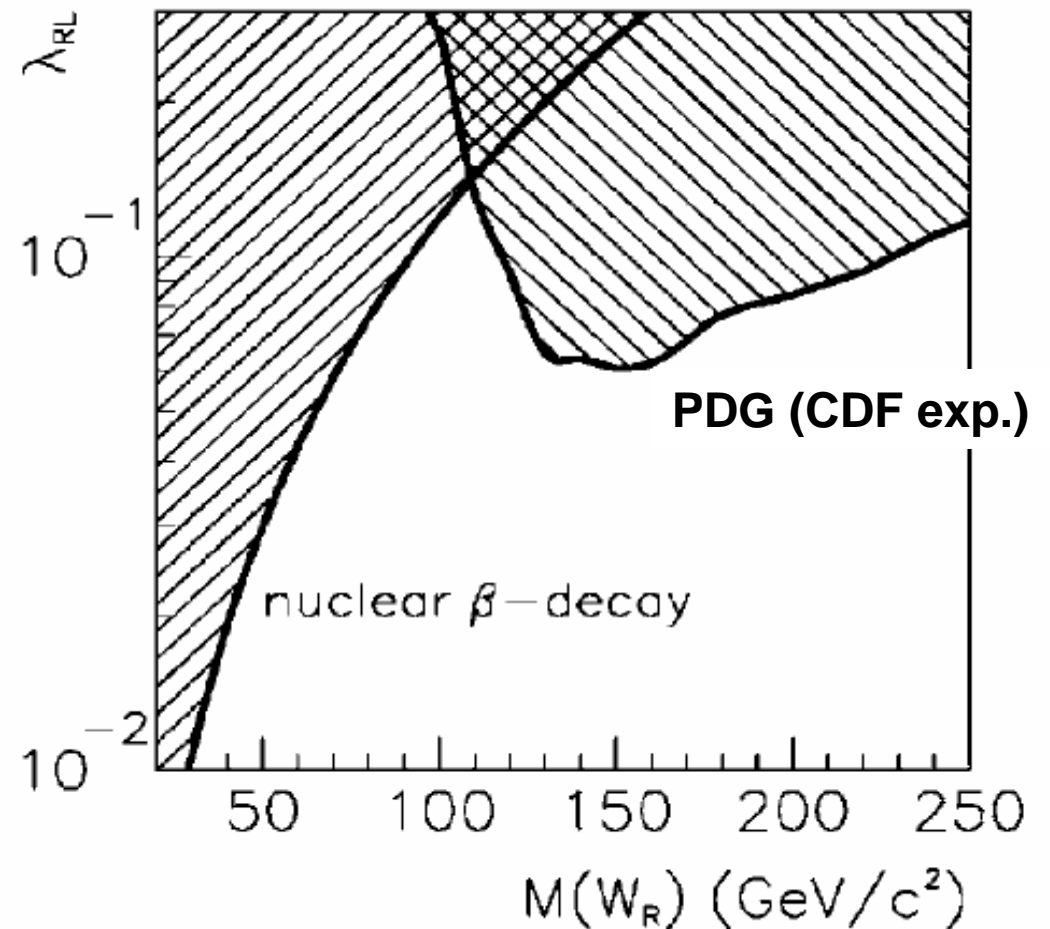
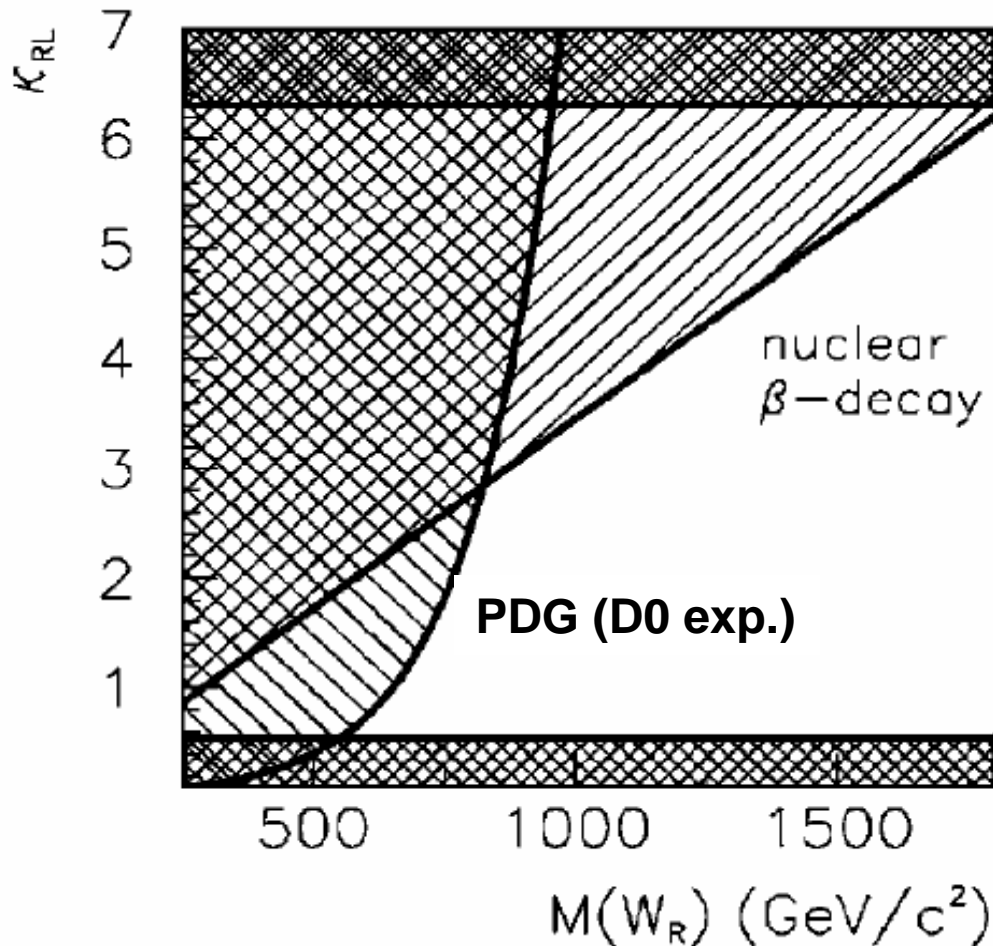
B. Time reversal

- 1. *D triple correlation***
- 2. *R triple correlation***

this is for another talk



Complimentarity of beta decay RHC results and collider results, in general LRS models



$$K_{RL} = g_R / g_L$$

shaded areas are excluded

$$\lambda_{RL} = V_{ud}^R / V_{ud}^L$$